# 5.2 Curl, Line Integrals, and Circulation

# 5.2.1 Line Integrals

Consider a curve *C* in a Euclidean vector space  $\mathbb{R}^3$ . Let r(t) = [x(t), y(t), z(t)] be a parametric representation of *C*. Furthermore, let

 $F : \mathbb{R}^3 \to \mathbb{R}^3$  be a vector-valued function of r and let r'(t) be the tangent vector. The **line integral** is

$$\int_{C} \boldsymbol{F}(\boldsymbol{r}(t)) \cdot \boldsymbol{r}'(t) \,\mathrm{d}t \tag{5.8}$$

which integrates *F* along the curve. For more on computing line integrals, see (Schey 2005; pp. 63-74) and (Kreyszig 2011; § 10.1 and 10.2).

If *F* is a **force** being applied to an object moving along the curve *C*, the line integral is the **work** done by the force. More generally, the line integral integrates *F* along the tangent of *C*.

# 5.2.2 Circulation

Consider the line integral over a closed curve *C*, denoted by

$$\oint_{C} F(\mathbf{r}(t)) \cdot \mathbf{r}'(t) \,\mathrm{d}t. \tag{5.9}$$

We call this quantity the **circulation** of *F* around *C*.

For certain vector-valued functions *F*, the circulation is zero for every curve. In these cases (static electric fields, for instance), this is sometimes called the **the law of circulation**.

# 5.2.3 Curl

Consider the division of the circulation around a curve in  $\mathbb{R}^3$  by the surface area it encloses  $\Delta S$ ,

$$\frac{1}{\Delta S} \oint_{C} F(\mathbf{r}(t)) \cdot \mathbf{r}'(t) \,\mathrm{d}t.$$
(5.10)

In a manner analogous to the operation that gaves us the divergence, let's consider shrinking this curve to a point and the surface area to zero,

$$\lim_{\Delta S \to 0} \frac{1}{\Delta S} \oint_{C} F(\mathbf{r}(t)) \cdot \mathbf{r}'(t) \,\mathrm{d}t.$$
(5.11)

We call this quantity the "scalar" **curl** of *F* at each point in  $\mathbb{R}^3$  *in the direction normal to*  $\Delta S$  as it shrinks to zero. Taking three (or *n* for  $\mathbb{R}^n$ ) "scalar" curls in independent



normal directions (enough to span the vector space), we obtain the **curl** proper, which is a vector-valued function curl :  $\mathbb{R}^3 \to \mathbb{R}^3$ .

The curl is coordinate-independent. In cartesian coordinates, it can be shown to be equivalent to the following.

Equation 5.12 curl: differential form, cartesian coordinates

$$\operatorname{curl} \boldsymbol{F} = \begin{bmatrix} \partial_{y} F_{z} - \partial_{z} F_{y} & \partial_{z} F_{x} - \partial_{x} F_{z} & \partial_{x} F_{y} - \partial_{y} F_{x} \end{bmatrix}^{\mathsf{T}}$$

But what does the curl of *F* represent? It quantifies the local rotation of *F* about each point. If *F* represents a fluid's velocity, curl *F* is the local rotation of the fluid about each point and it is called the **vorticity**.

# 5.2.4 Zero Curl, Circulation, and Path Independence

**5.2.4.1 Circulation** It can be shown that if the circulation of *F* on all curves is zero, then in each direction *n* and at every point curl F = 0 (i.e.  $n \cdot \text{curl } F = 0$ ). Conversely, for curl F = 0 in a simply connected region<sup>2</sup>, *F* has zero circulation.

Succinctly, informally, and without the requisite qualifiers above,

zero circulation 
$$\Rightarrow$$
 zero curl (5.13)

zero curl + simply connected region 
$$\Rightarrow$$
 zero circulation. (5.14)

**5.2.4.2 Path Independence** It can be shown that if the path integral of *F* on all curves between any two points is **path-independent**, then in each direction *n* and at every point curl F = 0 (i.e.  $n \cdot \text{curl } F = 0$ ). Conversely, for curl F = 0 in a simply connected region, all line integrals are independent of path.

Succinctly, informally, and without the requisite qualifiers above,

path independence  $\Rightarrow$  zero curl (5.15)

zero curl + simply connected region  $\Rightarrow$  path independence. (5.16)

# 5.2.4.3 And How They Relate It is also true that

path independence  $\Leftrightarrow$  zero circulation. (5.17)

So, putting it all together, we get figure 5.5.

2. A region is simply connected if every curve in it can shrink to a point without leaving the region. An example of a region that is not simply connected is the surface of a toroid.



Figure 5.5. An implication graph relating zero curl, zero circulation, path independence, and connectedness. Blue edges represent implication (*a* implies *b*) and black edges represent logical conjunctions.

#### 5.2.5 Exploring Curl

Curl is perhaps best explored by considering it for a vector field in  $\mathbb{R}^2$ . Such a field in cartesian coordinates  $F = F_x \hat{i} + F_y \hat{j}$  has curl

$$\operatorname{curl} \boldsymbol{F} = \begin{bmatrix} \partial_y 0 - \partial_z F_y & \partial_z F_x - \partial_x 0 & \partial_x F_y - \partial_y F_x \end{bmatrix}^{\top}$$
$$= \begin{bmatrix} 0 - 0 & 0 - 0 & \partial_x F_y - \partial_y F_x \end{bmatrix}^{\top}$$
$$= \begin{bmatrix} 0 & 0 & \partial_x F_y - \partial_y F_x \end{bmatrix}^{\top}.$$
(5.18)

That is,  $\operatorname{curl} \mathbf{F} = (\partial_x F_y - \partial_y F_x)\hat{\mathbf{k}}$  and the only nonzero component is normal to the *xy*-plane. If we overlay a quiver plot of  $\mathbf{F}$  over a "color density" plot representing the  $\hat{\mathbf{k}}$ -component of curl  $\mathbf{F}$ , we can increase our intuition about the curl. First, load some Python packages.

```
import numpy as np
import sympy as sp
import matplotlib.pyplot as plt
from matplotlib.ticker import LogLocator
from matplotlib.colors import *
```

Now we define some symbolic variables and functions.

```
x = sp.Symbol('x', real=True)
y = sp.Symbol('y', real=True)
F_x = sp.Function('F_x')(x, y)
F_y = sp.Function('F_y')(x, y)
```

We use a variation of the quiver\_plotter\_2D() from above to make several of these plots.

```
def quiver_plotter_2D(
  field={F_x: x*y, F_y: x*y},
  grid_width=3,
  grid_decimate_x=8,
  grid_decimate_y=8,
  norm=Normalize(),
  density_operation='div',
 print_density=True,
):
  # Define symbolics
  x, y = sp.symbols('x y', real=True)
  F_x = sp.Function('F_x')(x, y)
  F_y = sp.Function('F_y')(x, y)
  field_sub = field
  # Compute density
  if density_operation == 'div':
    den = F_x.diff(x) + F_y.diff(y)
  elif density_operation == 'curl':
    den = F_y.diff(x) - F_x.diff(y) # in the k direction
  else:
    raise ValueError('div and curl are the only density operators')
  den_simp = den.subs(field_sub).doit().simplify()
  if den_simp.is_constant():
    print('Warning: density operator is constant (no density plot)')
  if print_density:
    print(f'The {density_operation} is: {den_simp}')
  # Lambdify for numerics
  F_x_{sub} = F_x_{subs}(field_{sub})
  F_y_sub = F_y.subs(field_sub)
  F_x_fun = sp.lambdify((x, y),F_x.subs(field_sub), 'numpy')
  F_y_fun = sp.lambdify((x, y), F_y.subs(field_sub), 'numpy')
  if F_x_sub.is_constant:
    F_x_fun1 = F_x_fun # Dummy
    F_x_fun = lambda x, y: F_x_fun1(x, y)*np.ones(x.shape)
  if F_y_sub.is_constant:
    F_y_fun1 = F_y_fun # Dummy
    F_y_{fun} = lambda x, y: F_y_{fun1}(x, y) * np.ones(x.shape)
  if not den_simp.is_constant():
    den_fun = sp.lambdify((x, y), den_simp, 'numpy')
  # Create grid
  w = grid_width
  Y, X = np.mgrid[-w:w:100j, -w:w:100j]
  # Evaluate numerically
  F_x_num = F_x_fun(X, Y)
  F_y_num = F_y_fun(X, Y)
  if not den_simp.is_constant():
    den_num = den_fun(X, Y)
```

Let's inspect several cases.

```
fig, ax = quiver_plotter_2D(
   field={F_x: 0, F_y: sp.cos(2*sp.pi*x)}, density_operation='curl',
   grid_decimate_x=2, grid_decimate_y=10, grid_width=1
)
```

plt.draw()









# 5.3 Gradient

# 5.3.1 Gradient

The **gradient** grad of a scalar-valued function  $f : \mathbb{R}^3 \to \mathbb{R}$  is a vector field  $F : \mathbb{R}^3 \to \mathbb{R}^3$ ; that is, grad *f* is a vector-valued function on  $\mathbb{R}^3$ . The

gradient's local **direction** and **magnitude** are those of the local maximum rate of increase of *f*. This makes it useful in optimization (e.g., in the method of gradient descent).

This principle tells us that nature's laws quite frequently seem to be derivable by assuming a certain quantity—called *action*—is minimized. Considering, then, that the gradient supplies us with a tool for optimizing functions, it is unsurprising that the gradient enters into the equations of motion of many physical quantities.

The gradient is coordinate-independent, but its coordinate-free definitions don't add much to our intuition.

# Equation 5.19 gradient: cartesian coordinates

grad 
$$f = \begin{bmatrix} \partial_x f & \partial_y f & \partial_z f \end{bmatrix}^\top$$

# 5.3.2 Vector Fields from Gradients Are Special

Although all gradients are vector fields, not all vector fields are gradients. That is, given a vector field F, it may or may not be equal to the gradient of any scalarvalued function f. Let's say of a vector field that is a gradient that it has **gradience**.<sup>3</sup> Those vector fields that *are* gradients have special properties. Surprisingly, those properties are connected to path independence and curl. It can be shown that iff a field is a gradient, line integrals of the field are path independent. That is, for a vector field,

gradience 
$$\Leftrightarrow$$
 path independence. (5.20)

Considering what we know from section 5.2 about path independence we can expand figure 5.5 to obtain figure 5.10.

3. This is nonstandard terminology, but we're bold.

