7.2 Sturm-Liouville Problems

Before we introduce an important solution method for PDEs in section 7.3, we consider an *ordinary* differential equation that will

arise in that method when dealing with a single spatial dimension *x*: the **sturmliouville (S-L) differential equation**. Let *p*, *q*, σ be functions of *x* on open interval (*a*, *b*). Let *X* be the dependent variable and λ constant. The **regular S-L problem** is the S-L ODE²

$$\frac{\mathrm{d}}{\mathrm{d}x}\left(pX'\right) + qX + \lambda\sigma X = 0 \tag{7.1}$$

with boundary conditions

$$\beta_1 X(a) + \beta_2 X'(a) = 0 \tag{7.2}$$

$$\beta_3 X(b) + \beta_4 X'(b) = 0 \tag{7.3}$$

with coefficients $\beta_i \in \mathbb{R}$. This is a type of **boundary value problem**.

This problem has nontrivial solutions, called **eigenfunctions** $X_n(x)$ with $n \in \mathbb{Z}_+$, corresponding to specific values of $\lambda = \lambda_n$ called **eigenvalues**.³ There are several important theorems proven about this (see (Haberman 2018; § 5.3)). Of greatest interest to us are that

- 1. there exist an infinite number of eigenfunctions *X_n* (unique within a multiplicative constant),
- 2. there exists a unique corresponding *real* eigenvalue λ_n for each eigenfunction X_n ,
- 3. the eigenvalues can be ordered as $\lambda_1 < \lambda_2 < \cdots$,
- 4. eigenfunction X_n has n 1 zeros on open interval (a, b),
- 5. the eigenfunctions X_n form an orthogonal basis with respect to weighting function σ such that any piecewise continuous function $f : [a, b] \to \mathbb{R}$ can be represented by a generalized fourier series on [a, b].

This last theorem will be of particular interest in section 7.3.



^{2.} For the S-L problem to be *regular*, it has the additional constraints that p, q, σ are continuous and p, $\sigma > 0$ on [a, b]. This is also sometimes called the sturm-liouville eigenvalue problem. See (Haberman 2018; § 5.3) for the more general (non-regular) S-L problem and (§ 7.4) for the multi-dimensional analog. 3. These eigenvalues are closely related to, but distinct from, the "eigenvalues" that arise in systems of linear ODEs.

7.2.1 Types of Boundary Conditions

Boundary conditions of the sturm-liouville kind equation (7.2) have four sub-types:

dirichlet for just β_2 , $\beta_4 = 0$, **neumann** for just β_1 , $\beta_3 = 0$, **robin** for all $\beta_i \neq 0$, and **mixed** if $\beta_1 = 0$, $\beta_3 \neq 0$; if $\beta_2 = 0$, $\beta_4 \neq 0$.

There are many problems that are *not* regular sturm-liouville problems. For instance, the right-hand sides of equation (7.2) are zero, making them **homogeneous boundary conditions**; however, these can also be nonzero. Another case is **periodic boundary conditions**:

$$X(a) = X(b) \tag{7.4}$$

$$X'(a) = X'(b).$$
 (7.5)

Example 7.1

Consider the differential equation

$$X'' + \lambda X = 0$$

with dirichlet boundary conditions on the boundary of the interval [0, L]

X(0) = 0 and X(L) = 0.

Solve for the eigenvalues and eigenfunctions.

This is a sturm-liouville problem, so we know the eigenvalues are real. The well-known general solutions to the ODE is

$$X(x) = \begin{cases} k_1 + k_2 x & \lambda = 0\\ k_1 e^{j\sqrt{\lambda}x} + k_2 e^{-j\sqrt{\lambda}x} & \text{otherwise} \end{cases}$$

with real constants k_1 , k_2 . The solution must also satisfy the boundary conditions. Let's apply them to the case of $\lambda = 0$ first:

$$X(0) = 0 \Longrightarrow k_1 + k_2(0) = 0 \Longrightarrow k_1 = 0$$
$$X(L) = 0 \Longrightarrow k_1 + k_2(L) = 0 \Longrightarrow k_2 = -k_1/L$$

Together, these imply $k_1 = k_2 = 0$, which gives the *trivial solution* X(x) = 0, in which we aren't interested. We say, then, for nontrivial solutions $\lambda \neq 0$. Now let's check $\lambda < 0$. The solution becomes

$$X(x) = k_1 e^{-\sqrt{|\lambda|}x} + k_2 e^{\sqrt{|\lambda|}x}$$
$$= k_3 \cosh(\sqrt{|\lambda|}x) + k_4 \sinh(\sqrt{|\lambda|}x)$$

where k_3 and k_4 are real constants. Again applying the boundary conditions:

$$\begin{aligned} X(0) &= 0 \Longrightarrow k_3 \cosh(0) + k_4 \sinh(0) = 0 \Longrightarrow k_3 + 0 = 0 \Longrightarrow k_3 = 0 \\ X(L) &= 0 \Longrightarrow 0 \cosh(\sqrt{|\lambda|}L) + k_4 \sinh(\sqrt{|\lambda|}L) = 0 \Longrightarrow k_4 \sinh(\sqrt{|\lambda|}L) = 0. \end{aligned}$$

However, $\sinh(\sqrt{|\lambda|}L) \neq 0$ for L > 0, so $k_4 = k_3 = 0$ —again, the trivial solution. Now let's try $\lambda > 0$. The solution can be written

$$X(x) = k_5 \cos(\sqrt{\lambda}x) + k_6 \sin(\sqrt{\lambda}x).$$

Applying the boundary conditions for this case:

$$X(0) = 0 \Longrightarrow k_5 \cos(0) + k_6 \sin(0) = 0 \Longrightarrow k_5 + 0 = 0 \Longrightarrow k_5 = 0$$
$$X(L) = 0 \Longrightarrow 0 \cos(\sqrt{\lambda}L) + k_6 \sin(\sqrt{\lambda}L) = 0 \Longrightarrow k_6 \sin(\sqrt{\lambda}L) = 0.$$

Now, $\sin(\sqrt{\lambda}L) = 0$ for

$$\sqrt{\lambda}L = n\pi \Longrightarrow$$
$$\lambda = \left(\frac{n\pi}{L}\right)^2. \qquad (n \in \mathbb{Z}_+)$$

Therefore, the only nontrivial solutions that satisfy both the ODE and the boundary conditions are the *eigenfunctions*

$$X_n(x) = \sin\left(\sqrt{\lambda_n}x\right) \tag{7.6}$$

$$=\sin\left(\frac{n\pi}{L}x\right) \tag{7.7}$$

with corresponding *eigenvalues*

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2.$$

Note that because $\lambda > 0$, λ_1 is the lowest eigenvalue.

Plotting the Eigenfunctions

```
import numpy as np
import matplotlib.pyplot as plt
```

Set L = 1 and compute values for the first four eigenvalues lambda_n and eigenfunctions X_n.



We see that the fourth of the S-L theorems appears true: n - 1 zeros of X_n exist on the open interval (0, 1).

7.3 PDE Solution by Separation of Variables

We are now ready to learn one of the most important techniques for solving PDEs: **separation of variables**. It applies only to **linear** PDEs

since it will require the principle of superposition. Not all linear PDEs yield to this solution technique, but several that are important do.

The technique includes the following steps.

- **assume a product solution** Assume the solution can be written as a **product solution** *u_v*: the product of functions of each independent variable.
- **separate PDE** Substitute u_p into the PDE and rearrange such that at least one side of the equation has functions of a single independent variabe. If this is possible, the PDE is called **separable**.
- **set equal to a constant** Each side of the equation depends on different independent variables; therefore, they must each equal the same constant, often called $-\lambda$.
- **repeat separation**, **as needed** If there are more than two independent variables, there will be an ODE in the separated variable and a PDE (with one fewer variables) in the other independent variables. Attempt to separate the PDE until only ODEs remain.
- **solve each boundary value problem** Solve each boundary value problem ODE, ignoring the initial conditions for now.
- **solve the time variable ODE** Solve for the general solution of the time variable ODE, sans initial conditions.
- **construct the product solution** Multiply the solution in each variable to construct the product solution u_p . If the boundary value problems were sturm-liouville, the product solution is a family of **eigenfunctions** from which any function can be constructed via a generalized fourier series.
- **apply the initial condition** The product solutions individually usually do not meet the initial condition. However, a generalized fourier series of them nearly always does. **Superposition** tells us a linear combination of solutions to the PDE and boundary conditions is also a solution; the unique series that also satisfies the initial condition is the unique solution to the entire problem.

Example 7.2

Consider the one-dimensional diffusion equation PDE^a

$$\partial_t u(t, x) = k \partial_{xx}^2 u(t, x)$$

with real constant *k*, with dirichlet boundary conditions on inverval $x \in [0, L]$

$$u(t,0) = 0 \tag{7.8}$$

$$u(t, L) = 0,$$
 (7.9)



